

TURBULENT MIXING AND PARTICLE  
DISTRIBUTION INVESTIGATIONS  
IN THE HIMMERFJÄRDEN 1978

By Erik Buch



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fördelning i Himmerfjärden 1978

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## ABSTRACT

Turbulent diffusion in the weak halocline in the Himmerfjärd in the fall was studied by a tracer experiment. The tracer, Rhodamine B, was injected instantaneously at a prescribed depth and after that the spreading of the dye patch was followed by a systematic mapping of the concentration distribution.

Vertical exchange coefficients computed from the experiment range from  $1.9 \cdot 10^{-2}$  to  $3.1 \cdot 10^{-2}$  cm<sup>2</sup>/s. The vertical exchange is rather weak in comparison with results from open waters at the same windspeed and stratification. This is thought to be due to a reduction of the wind induced mixing caused by the physical boundaries of the narrow fiord.

Horizontal diffusion characteristics are computed from the time rate of change of the concentration distribution variances. Apparent horizontal diffusivities range up to about  $5 \cdot 10^3$  cm<sup>2</sup>/s when the observed scale of turbulence ranges up to about 800 m.

The results are not consistent with a model for generation of horizontal dispersion by vertical current shear and vertical diffusion and it is concluded that horizontal diffusion and also horizontal shear are important factors for generation of dispersion in the Himmerfjärd.

Particle distribution in the Himmerfjärd was investigated by measuring vertical profiles of light attenuation at a number of stations. The results show a layering of the fiord into three distinct layers such that the particle content is relatively high in the upper and the bottom layer and relatively low in the intermediate layer. The intermediate layer is believed to be intruding Baltic water.

The work has been carried out by Dr Erik Buch, Institute of Physical Oceanography, University of Copenhagen. It is a part of the project "Materialtransporter i Himmerfjärden" financed by grants to SMHI from the Research Department of the Swedish Environment Protection Board, contract nr 7-329/79. This project is intended to give hydrographic information for the system ecological study "Himmerfjärdsundersökningen" run by the Askö Laboratory. Project leader at SMHI is Erland Bergstrand. The tracer study was initiated in a cooperation between Dr Wayne Wilmot at the RGASA group, University of Stockholm and Dr Erik Buch.

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INTRODUCTION

The Himmerfjärd has been thoroughly investigated physically, chemically and biologically during the last years in order to develop models describing the processes which are judged to be important for longterm balances of material.

For the modelling there is a need of information about the horizontal and vertical turbulent diffusion and the particle distribution in the Himmerfjärd, wherefore a number of dye diffusion- and optical light transmission experiments has been carried out during the week 2 - 6 October 1978 onboard the research vessel "Aurelia".

INSTRUMENTATION AND EXPERIMENTAL TECHNIQUE

1. Diffusion experiments

The turbulent diffusion has been investigated by means of a dye tracing method using the organic, fluorescent liquid Rhodamine B as a tracer. An in situ fluorometer, described by Kullenberg (1968, 1974), has been used for the tracing. Information about the current, especially the current shear, and the density stratification during the period of diffusion experiments was gained by means of a permanent current meter station and thermistor chain station measuring the current, the salinity and the temperature every two meters in the depth interval 4 - 16 meters. The dye experiments were carried out in the pycnocline layer, because it is of special interest to know the mixing between the two layers in stratified waters.

The density of the dye was adjusted to the density of the depth of injection within  $0.1 \sigma_t$  - units, and the dye was injected by

means of high pressure of 4 - 10 atm through a hose connected to a two meter long diffusor. The injection lasted for about five minutes which means that the release of dye may be regarded as momentaneous after 2 - 3 hours, according to the criterion given by Okubo (1971). When the dye patch, mainly because of the current shear, was spread out horizontally a systematic mapping of the concentration distribution started.

The fluorometer, in tow at the shipside, was cycling in a depth interval around the depth of injection. During the experiments in the Himmerfjärd radar bearings were used for position determinations.

The duration of the experiment was 12 - 16 hours. Observations of the wind were made at intervals during the whole experiment.

## 2. Particle distribution investigations

For determination of the particle distribution a light attenuation meter (c - meter) was used. The path length of the c - meter is 1.4 m and it measures the light attenuation coefficient  $c$  in the red ( $\lambda_{\max} = 655 \text{ nm}$ ), the green ( $\lambda_{\max} = 525 \text{ nm}$ ) and the ultraviolet ( $\lambda_{\max} = 380 \text{ nm}$ ) part of the spectrum. However, in the Himmerfjärd the content of material, especially yellow substance, was so great that all ultraviolet light was totally attenuated along the 1.4 m long path.

Vertical profiles were made at 11 stations placed in such a way that longitudinal and transversal sections of particle distribution were obtained, see Fig 1.

## RESULTS

Symbols used are listed on page 15.

1. Vertical diffusion

Assuming that

- a) the advective terms may be disregarded because the observations are made in a Lagrangian reference system with origo at the center of the dye patch,
- b) the vertical velocity is negligible,
- c) the decrease in the dye concentration is due only to vertical diffusion, and
- d) the vertical diffusion coefficient  $K$  is a constant, the general diffusion equation is reduced to

$$1) \quad \frac{\partial C}{\partial t} = K \cdot \frac{\partial^2 C}{\partial z^2}$$

Because radar navigation was used in the Himmerfjärd, the observations were made in an absolute reference system, but during the data processing they were transformed into a Lagrangian system by considering the drift velocities, so that the assumption a) still holds.

A solution to eq 1 has been given by Kullenberg (1968):

$$2) \quad K = \frac{H^2}{\pi (t_2 - t_1)} \ln \frac{C_1}{C_2}$$

where subscripts 1 and 2 denote successive times of observation. Identification of a given layer with increasing diffusion time has been done by means of thickness, shape and position in the temperature structure.

The mixing coefficients obtained by the experiments in the Himmerfjärd are given in Table 2.

Kullenberg (1971; 1974) found theoretically the relation



$$3) \quad K = Rf \ c_d \ \frac{\rho_a}{\rho_o} \left( \frac{W_a}{N} \right)^2 \left| \frac{dq}{dz} \right|$$

by assuming that

- a) the turbulence is shear generated, and
- b) the internal stress is approximately equal to the surface wind stress.

The coefficient of proportionality ( $Rf \ c_d \ \frac{\rho_a}{\rho_o}$ ) can be found from the observations by plotting  $K$  versus  $\left( \frac{W_a}{N} \right)^2 \left| \frac{dq}{dz} \right|$

From investigations in open sea- and coastal areas, Kullenberg found:

$$4) \quad Rf \ c_d \ \frac{\rho_a}{\rho_o} = 9.6 \cdot 10^{-8}$$

for persisting winds ( $W_a \geq 4-5$  m/s).

In Fig 17 the results from the present investigations in the Himmerfjärd are given together with results obtained by the author in areas similar to the Himmerfjärd, Buch (1978, 1979). It is seen that the coefficient

$$5) \quad Rf \ c_d \ \frac{\rho_a}{\rho_o} = 1.3 \cdot 10^{-9}$$

is considerably smaller than the value obtained by Kullenberg. Mixing experiments giving this low vertical mixing have all been carried out in narrow fiords and so it is believed that the physical boundaries are reducing the mixing effect due to the wind.

In Table 4 is given an evaluation of the entrainment from the lower to the upper layer. The entrainment velocity,  $u_e$ , is here calculated as

$$u_e = \frac{KN}{g'} \quad (\text{Kullenberg 1977})$$

In a theory of Kato and Phillips (1969) for entrainment due to shear generated turbulence there is given the relation

$$u_e / u_* = 2.5 \text{ Ri}_*^{-1}$$

In Table 4 the quantities of this relation are calculated whereby the friction velocity  $u_*$  is assumed to be given by the wind only. The reduced wind fetch mentioned above will here have the effect that the constant of proportionality between  $u_e / u_*$  and  $\text{Ri}_*^{-1}$  is very much lower than the one obtained by Kato and Phillips (1969).

## 2. Horizontal diffusion

A mean picture of the horizontal dye concentration distribution has been obtained by plotting the observed concentrations along the track of the ship and contouring the plot. To get a picture as accurate as possible the dye patch must be covered in a number of directions but, on the other hand, the time used for one proper covering of the patch must not exceed more than about 20% of the diffusion time.

Examples of concentration distribution pictures obtained in the Himmerfjärd are shown in Figs 7 - 16.

To get information about the horizontal diffusion characteristics from the dye distribution the theory of Joseph and Sendner (1958) can be used. The observed distribution of the dye has been transformed into a rotationally symmetric one by calculating the equivalent radius  $r$ , i.e. the radius of a circle with an area equivalent to the area inside a given concentration isoline. The theory gives an equation for calculation of the horizontal diffusion velocity  $P$ :

$$6) \quad \sigma_{rc}^2 = 6 P^2 t^2$$

where

$$7) \quad \sigma_{rc}^2 = \frac{\int_0^{\infty} r^2 C(r,t) 2\pi r dr}{\int_0^{\infty} C(r,t) 2\pi r dr}$$

is the definition of the rotationally symmetric variances, see Table 5.

As seen in Figs 7 - 16 the distribution of dye is not rotationally symmetric, but is more like an ellipse. Therefore it may be more accurate to find the longitudinal and transversal variances

$$\sigma_x^2 \text{ and } \sigma_y^2$$

defined by

$$8) \quad \sigma_x^2 = \frac{\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} x^2 C dx dy}{\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} C dx dy}$$

$$9) \quad \sigma_y^2 = \frac{\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} y^2 C dx dy}{\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} C dx dy}$$

The vertical distribution is assumed homogenous.

If the dye distribution is Gaussian, the definition of  $\sigma_{rc}^2$  writes

$$10) \quad \sigma_{rc}^2 = \sigma_x \sigma_y$$

In Table 6 a comparison of  $\sigma_{rc}^2$  and  $\sigma_x \sigma_y$  is given. As expected, there is some scattering in the results, owing partly to experimental uncertainties and partly to the fact that the dye distribution is neither rotationally symmetric nor elliptic.

The rotationally symmetric variance  $\sigma_{rc}^2$  is also used for definition of the apparent horizontal diffusivity.

$$11) \quad K_h = \frac{\sigma_{rc}^2}{4t}$$

and the scale of turbulence

$$12) \quad l_h = 3\sigma_{rc}$$

By plotting  $\sigma_{rc}^2$  versus the time  $t$  and  $K_h$  versus  $l_h$  the following relations are found

$$13) \quad \sigma_{rc}^2 \sim t^{2.2}$$

$$14) \quad K_h \sim l_h^{4/3}$$

As seen from Figs 18 and 19 the present results are very similar to the results of Kullenberg (1974), which means that the dye patch has spread out horizontally in a way that could be foreseen.  $\sigma_{rc}^2 \sim t^{2.2}$  would by using eq 11) and 12) give  $K_h \sim l_h^{1.1}$ . This inconsistency with the observed result may be due to the uncertainties in determining  $\sigma_{rc}$ .

In many areas it is found that the horizontal dispersion is determined by the vertical diffusion and the vertical current shear, and often to such a degree that the horizontal diffusion may be neglected, as exemplified by the findings of Kullenberg (1968).

$$15) \quad P = (1.1 H_o \frac{dq}{dz} + 1.3) \times 10^{-2}$$

By considering the shear - diffusion effect in an internal unbounded layer and by neglecting the vertical convection, the horizontal diffusion and the horizontal shear, Kullenberg (1971, 1974) developed a model which describes the apparent horizontal dispersion generated by the vertical diffusion and the vertical current shear.

In case of an instantaneous cubic source of dye with an initial thickness  $H_o$ , the model leads to a prediction of the horizontal variances:

$$16) \quad \sigma_x^2 = (2/3 a_o^2 t^3 + \frac{a^2}{\omega^2} t) K + \frac{1}{12} a_o^2 H_o^2 t^2$$

$$17) \quad \sigma_y^2 = \left( 9/8 \frac{b^2}{\omega^2} t + \frac{b^2 a}{a^2 \omega^2} (2 a_o - a) \right) K$$

The results obtained using this model are given in Table 7. Comparing the theoretical horizontal variances with those obtained from the experiments given in Table 6, it appears that the theoretical values are much lower than the observed ones. This implies that some of the assumptions made in developing the model do not hold in the Himmerfjärd and similar areas, Buch (1978, 1979). If the apparent diffusion velocity P is calculated from eq. 15, see Table 7, values about an order of magnitude lower than the observed ones is found, which indicates that the horizontal diffusion may not be neglected. The horizontal shear is also expected to be important in narrow areas like the Himmerfjärd.

### 3. Particle distribution

The attenuation coefficient  $c$  at the wavelength  $\lambda$  and the depth  $z$  consists of three components:

$$18) \quad c(\lambda, z) = c_p(\lambda, z) + a_y(\lambda, z) + c_w(\lambda)$$

The attenuation due to dissolved organic matter (yellow substance) is very wavelength selective with high attenuation in the ultraviolet and almost no attenuation in the green and red part of the spectrum. This has been utilized for an in situ determination of  $a_y$ , Jerlov (1976):

$$19) \quad a_y(380) = \left( c(380) - c_w(380) \right) - K \left( c(655) - c_w(655) \right)$$

where  $K$  is a factor dependent upon the particle characteristics. As mentioned above, the content of yellow substance in the Himmerfjärd was so great that the ultraviolet light was totally attenuated, which means that  $a_y(380)$  can not be determined.

However, the wavelength selectivity of yellow substance leads to:

$$20) \quad c(\lambda, z) = c_p(\lambda, z) + c_w(\lambda)$$

for

$$\lambda = \begin{cases} 525 \text{ nm} \\ 655 \text{ nm} \end{cases}$$

According to Clarke and James (1939), the attenuation coefficients for pure water are:

$$c_w(525) = 0.041 \text{ m}^{-1}; \quad c_w(655) = 0.304 \text{ m}^{-1}$$

Optical measurements were carried out three times during the period 2 - 6 October 1978.

From linear regression analysis the following relationship has been obtained:

$$21) \quad c(525) = 1.24 c(655) - 0.34$$

On this background only longitudinal and transversal sections showing the attenuation coefficient in the red due to particles are represented in this report, see Figs 20 - 25.

The observations show that the water in the Himmerfjärd can be divided into three distinct layers, a homogeneous upper layer with a relatively high content of particles, an intermediate layer of clear water, and a bottom layer rich on particles.

The behaviour of the three layers during the week of observations indicates an intrusion of clear and warm Baltic water (see Figs 2 - 6, 20 - 21) in the depth interval 6 - 12 meters during the days before 2 October. In the following days this water was mixed up with water from the other two layers, making the intermediate layer less distinct.

The observations from the 6 October indicate a new intrusion of Baltic water.

Unfortunately, the current meters between 6 - 10 meters were out of order during the whole period of observations, but the current meter at 12 meter showed in-going currents in the days up to the

2 October and again at the 6 October.

The upper layer had a rather homogeneous distribution of particles probably due to wind generated turbulent mixing.

The thickness of the upper layer was about 6 meters at the 2 October increasing to about 10 - 12 meters at the 6 October.

The increase in thickness of the upper layer can not be used directly for an estimation of the vertical diffusion coefficient, because the particle content is not a conservative property (waste water discharge, biological activity etc.) and sedimentation of particles from the upper to the intermediate layer may also be of some importance. In connection to the light attenuation measurements Secchi disc observations were carried out too, and a Secchi disc depth of 3.8 meters was found at all the optical stations.

A Secchi disc depth of 3.8 meters means that the depth of the euphotic zone, defined as the depth at which 1% of the surface quanta irradiance is found, lies in the depth interval 6.5 - 8 meters, Højerslev (1978).

According to Jerlows (1976) optical classification, the Himmerfjärd then belongs to coastal water type no. 9.

## CONCLUSIONS

The mixing experiments carried out in the Himmerfjärd October 1978 have shown that in such a narrow area the physical boundaries affect the mixing characteristics.

A dependence of the vertical mixing on the stratification, the current shear and the wind has been proved, but the constant of proportionality is much lower than the one found in open areas.

The horizontal dispersion progressed in a manner that could be expected from the diffusion diagrams given in Figs 18 - 19, but a prediction of the dispersion using Kullenbergs advection - diffusion model has indicated that horizontal diffusion and

especially horizontal shear are of importance in the Himmerfjärd. On this background it would be advisable in future experiments to investigate the horizontal shear.

The optical observations have shown a layering of the Himmerfjärd waters with three distinct layers, where the intermediate layer is believed to be intruding Baltic Water.

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## LIST OF SYMBOLS

a, b	oscillating vertical current shear
$a_o$	mean vertical current shear
$a_y$	light attenuation coefficient due to dissolved organic substance
$c_d$	drag coefficient
C	concentration, mass per unit volume
c	light attenuation coefficient
$c_p$	light attenuation coefficient due to particle
$c_w$	light attenuation coefficient due to pure water
g	acceleration of gravity
$g'$	reduced gravity
$H_o$	initial thickness of the dye layer
2h, H	thickness of the dye layer
$h_1$	depth to interfacial layer
K	vertical turbulent transport coefficient for mass
$K_x, K_y$	turbulent transport coefficients in the x, y directions, respectively
$M_o$	mass of injected dye
N	Brunt-Väisälä frequency
P	horizontal diffusion velocity
q	horizontal current vector
r	equivalent radius
Rf	flux Richardson number
Ri	Richardson number
$Ri_* =$	$\frac{g' \cdot h_1}{u_*^2}$ ; overall Richardson number
S	salinity
T	temperature
t	time
u, v, w	velocities in the x, y, z directions, respectively
U, V	mean horizontal velocities
$u', v', w'$	fluctuating velocities
$u_*$	friction velocity
$u_e$	entrainment velocity
$W_a$	wind velocity
x, y, z	Cartesian coordinate system
$\epsilon$	dissipation of turbulent energy per unit mass and time
$\rho$	density
$\rho_a$	density of the air
$\tau$	stress
$\tau_o$	wind stress at the sea surface
$\omega$	frequency of an oscillating current
$\sigma_t$	$= (\rho_{S,T,0}^{-1}) 10^3$
$\sigma_x^2, \sigma_y^2$	variances of the concentration distribution in the x, y directions
$\sigma_{rc}^2$	horizontal symmetrical variance

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Table 1. Environmental observations: Mean stability parameter  $N^2$ , Vertical shear, local Richardson number and wind  $W_a$ .

Date	Depth interval m	$N^2$ sec <sup>-2</sup>	$\frac{dq}{dz}$ sec <sup>-2</sup>	Ri	Wind m · sec <sup>-1</sup>
781002	12.5-15	$7.8 \cdot 10^{-4}$	$6.8 \cdot 10^{-3}$	16.8	7 - 8
781003	12.5-15	$8.2 \cdot 10^{-4}$	$1.5 \cdot 10^{-2}$	3.6	7 - 10
781005	12.5-15	$5.3 \cdot 10^{-4}$	$1.3 \cdot 10^{-2}$	3.1	6

Table 2. Thickness  $2h$ , time interval  $t_2 - t_1$ , concentration ratio  $C_1/C_2$ , mean depth, Vertical exchange coefficient  $K$  and mean  $K$ .

Date	2h cm	$t_2 - t_1$ sec	$C_1/C_2$	mean depth	$K$ cm <sup>2</sup> sec <sup>-1</sup>	mean $K$ cm <sup>2</sup> sec <sup>-1</sup>
781002	37	$2.7 \cdot 10^3$	2.00	12	$3.4 \cdot 10^{-2}$	
	36	$4.2 \cdot 10^3$	2.16	12	$2.4 \cdot 10^{-2}$	
	24	$1.6 \cdot 10^4$	22.09	12	$1.1 \cdot 10^{-2}$	$2.3 \cdot 10^{-2}$
781002	21	$1.1 \cdot 10^3$	4.20	12	$5.9 \cdot 10^{-2}$	
	20	$1.3 \cdot 10^4$	6.33	12	$6.5 \cdot 10^{-3}$	
	19	$3.5 \cdot 10^4$	20.00	12	$2.3 \cdot 10^{-3}$	$1.9 \cdot 10^{-2}$
781005	63	$1.5 \cdot 10^4$	4.06	12	$3.7 \cdot 10^{-2}$	
	68	$1.8 \cdot 10^4$	7.68	12	$4.8 \cdot 10^{-2}$	
	32	$1.5 \cdot 10^4$	6.37	12	$2.7 \cdot 10^{-2}$	
	32	$3.2 \cdot 10^3$	1.41	12	$1.2 \cdot 10^{-2}$	$3.1 \cdot 10^{-2}$

Table 3. Mean exchange coefficient  $\bar{K}$ , stability parameter  $\bar{N}^2$ , Shear  $\frac{dq}{dz}$ , square of the mean wind  $\bar{W}_a^2$  and the parameter  $\frac{W_a^2}{N^2} \frac{dq}{dz}$

Date	$\bar{K}$ cm <sup>2</sup> sec <sup>-1</sup>	$\bar{N}^2$ sec <sup>-2</sup>	$\frac{dq}{dz}$ sec <sup>-1</sup>	$\bar{W}_a^2$ cm <sup>2</sup> sec <sup>-2</sup>	$\frac{W_a^2}{N^2} \frac{dq}{dz}$ cm <sup>2</sup> sec <sup>-1</sup>
781002	$2.3 \cdot 10^{-2}$	$7.8 \cdot 10^{-4}$	$6.8 \cdot 10^{-3}$	$5.6 \cdot 10^5$	$4.9 \cdot 10^6$
781003	$1.9 \cdot 10^{-2}$	$8.2 \cdot 10^{-4}$	$1.5 \cdot 10^{-2}$	$6.1 \cdot 10^5$	$1.2 \cdot 10^7$
781005	$3.1 \cdot 10^{-2}$	$5.3 \cdot 10^{-4}$	$1.3 \cdot 10^{-2}$	$3.6 \cdot 10^5$	$8.8 \cdot 10^6$

Table 4. Mean vertical exchange coefficient  $\bar{K}$ , Stability parameter  $\bar{N}^2$ , reduced gravitational force, entrainment velocity, mean depth, square of the wind, overall Richardson number  $Ri_*$  and the ratio  $u_e/u_*$ .

Date	$\bar{K}$ cm <sup>2</sup> sec <sup>-1</sup>	$\bar{N}^2$ sec <sup>-1</sup>	$g \frac{\Delta \rho}{\rho}$ cm sec <sup>-2</sup>	$u_e$ cm sec <sup>-1</sup>	$h_1$ m	$\bar{W}_a^2$ cm <sup>2</sup> sec <sup>-2</sup>	$Ri_*$	$u_e/u_*$
781002	$2.3 \cdot 10^{-2}$	$7.8 \cdot 10^{-4}$	0.49	$3.7 \cdot 10^{-5}$	10	$5.6 \cdot 10^5$	$6.1 \cdot 10^2$	$4.1 \cdot 10^{-5}$
781003	$1.9 \cdot 10^{-2}$	$8.2 \cdot 10^{-4}$	0.59	$7.3 \cdot 10^{-6}$	13.5	$6.1 \cdot 10^5$	$6.1 \cdot 10^5$	$7.8 \cdot 10^{-6}$
781005	$3.1 \cdot 10^{-2}$	$5.3 \cdot 10^{-4}$	0.42	$3.9 \cdot 10^{-5}$	12.5	$3.6 \cdot 10^5$	$1.0 \cdot 10^3$	$5.4 \cdot 10^{-5}$

Table 5. Diffusion time, symmetrical variance  $\sigma_{rc}^2$ , Diffusion velocity P, apparent horizontal diffusion coefficient  $K_h$  and the horizontal lengthscale  $l_h$

Date	t sec	$\sigma_{rc}^2$ cm <sup>2</sup>	P cm sec <sup>-1</sup>	$K_h$ cm <sup>2</sup> sec <sup>-1</sup>	$l_h$ cm
781002	$3.6 \cdot 10^4$	$2.1 \cdot 10^8$	0.16	$1.5 \cdot 10^3$	$4.3 \cdot 10^4$
781003	$2.2 \cdot 10^4$	$3.1 \cdot 10^8$	0.33	$3.5 \cdot 10^3$	$5.3 \cdot 10^4$
	$2.7 \cdot 10^4$	$4.5 \cdot 10^8$	0.32	$4.2 \cdot 10^3$	$6.4 \cdot 10^4$
	$4.1 \cdot 10^4$	$6.3 \cdot 10^8$	0.25	$3.8 \cdot 10^3$	$7.5 \cdot 10^4$
781005	$2.4 \cdot 10^4$	$1.7 \cdot 10^8$	0.23	$1.8 \cdot 10^3$	$3.9 \cdot 10^4$
	$3.5 \cdot 10^4$	$2.1 \cdot 10^8$	0.17	$1.5 \cdot 10^3$	$4.3 \cdot 10^4$
	$4.6 \cdot 10^4$	$5.2 \cdot 10^8$	0.21	$2.8 \cdot 10^3$	$6.8 \cdot 10^4$

Table 6. Observed horizontal variances

Date	t sec	$\sigma_x^2$ cm <sup>2</sup>	$\sigma_y^2$ cm <sup>2</sup>	$\sigma_y/\sigma_x$	$\sigma_{rc}^2$ cm <sup>2</sup>	$\sigma_x \sigma_y$ cm <sup>2</sup>
781002	$3.6 \cdot 10^4$	$5.5 \cdot 10^8$	$4.7 \cdot 10^6$	0.27	$2.1 \cdot 10^8$	$1.5 \cdot 10^8$
781003	$2.7 \cdot 10^4$	$1.5 \cdot 10^9$	$1.3 \cdot 10^7$	0.29	$4.5 \cdot 10^8$	$4.4 \cdot 10^8$
	$4.1 \cdot 10^4$	$2.3 \cdot 10^9$	$3.4 \cdot 10^7$	0.38	$6.3 \cdot 10^8$	$9.0 \cdot 10^8$
781005	$3.5 \cdot 10^4$	$6.3 \cdot 10^8$	$2.2 \cdot 10^8$	0.59	$2.1 \cdot 10^8$	$3.7 \cdot 10^8$
	$4.6 \cdot 10^4$	$4.5 \cdot 10^9$	$1.3 \cdot 10^8$	0.17	$5.2 \cdot 10^8$	$7.7 \cdot 10^8$

Table 7. Parameters for prediction of variances and predicted variances

Date	t sec	$H_o$ cm	P (eq. 15)	K cm <sup>2</sup> sec <sup>-1</sup>	$a_o$ sec <sup>-1</sup>	a sec <sup>-1</sup>
781002	$3.6 \cdot 10^4$	200	0.03	$2.3 \cdot 10^{-2}$	$4.9 \cdot 10^{-3}$	$1.25 \cdot 10^{-3}$
781003	$4.1 \cdot 10^4$	200	0.05	$1.9 \cdot 10^{-2}$	$6.15 \cdot 10^{-3}$	$8.75 \cdot 10^{-3}$
781005	$3.5 \cdot 10^4$	200	0.04	$3.1 \cdot 10^{-3}$	$3.25 \cdot 10^{-3}$	$3.00 \cdot 10^{-3}$

Date	b sec <sup>-1</sup>	$\omega$ sec <sup>-1</sup>	$\sigma_x^2$ cm <sup>2</sup>	$\sigma_y^2$ cm <sup>2</sup>	$\sigma_x \sigma_y$ cm <sup>2</sup>	$\sigma_y/\sigma_x$
781002	$9.7 \cdot 10^{-3}$	$4.4 \cdot 10^{-4}$	$1.2 \cdot 10^8$	$1.0 \cdot 10^7$	$3.5 \cdot 10^7$	0.29
781003	$3.3 \cdot 10^{-3}$	$4.4 \cdot 10^{-4}$	$2.5 \cdot 10^8$	$1.8 \cdot 10^5$	$6.7 \cdot 10^6$	0.03
781005	$2.5 \cdot 10^{-3}$	$4.4 \cdot 10^{-4}$	$4.9 \cdot 10^7$	$8.5 \cdot 10^4$	$2.0 \cdot 10^6$	0.04

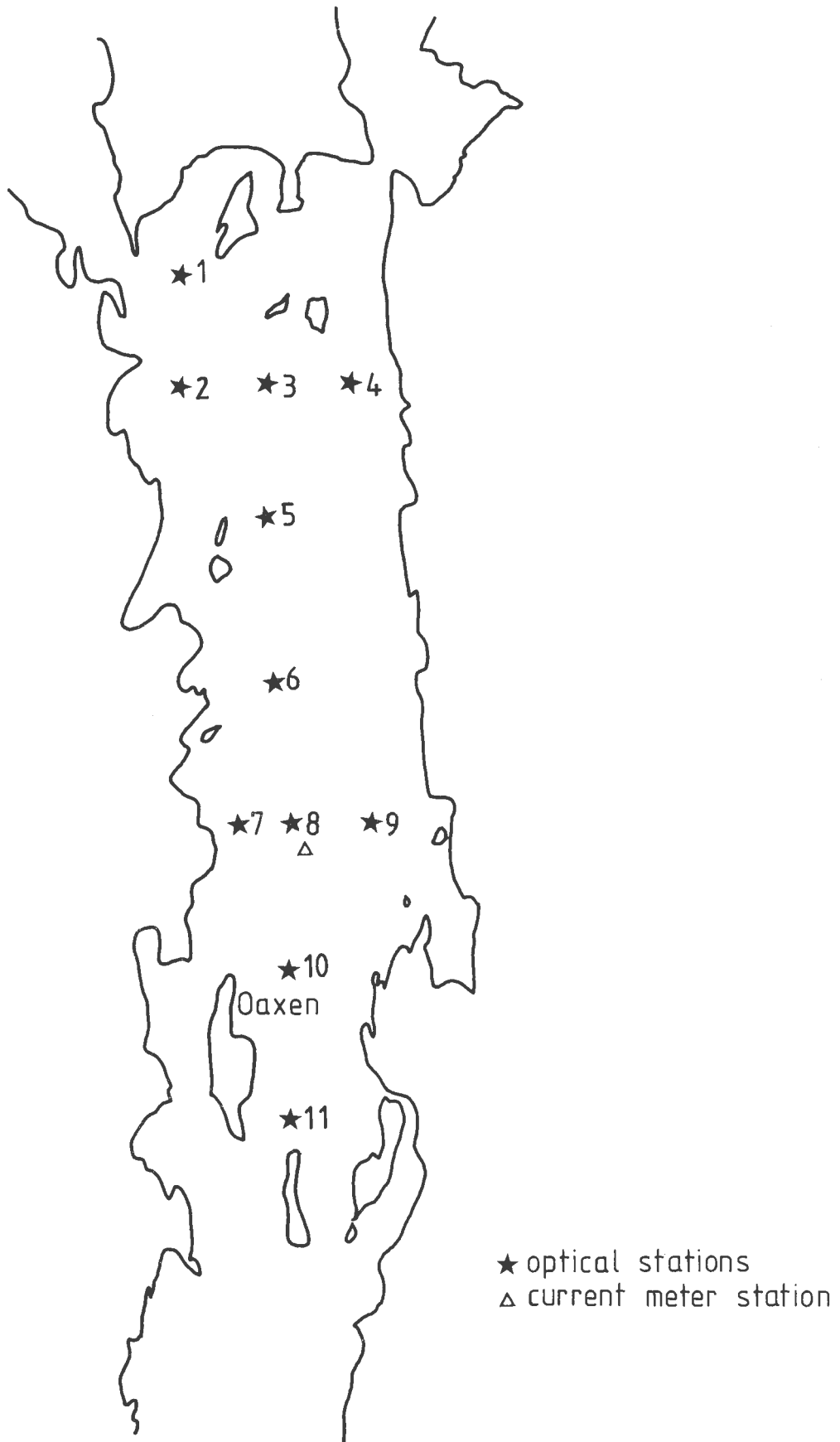


Figure 1. Area of investigation, the Himmerfjärd

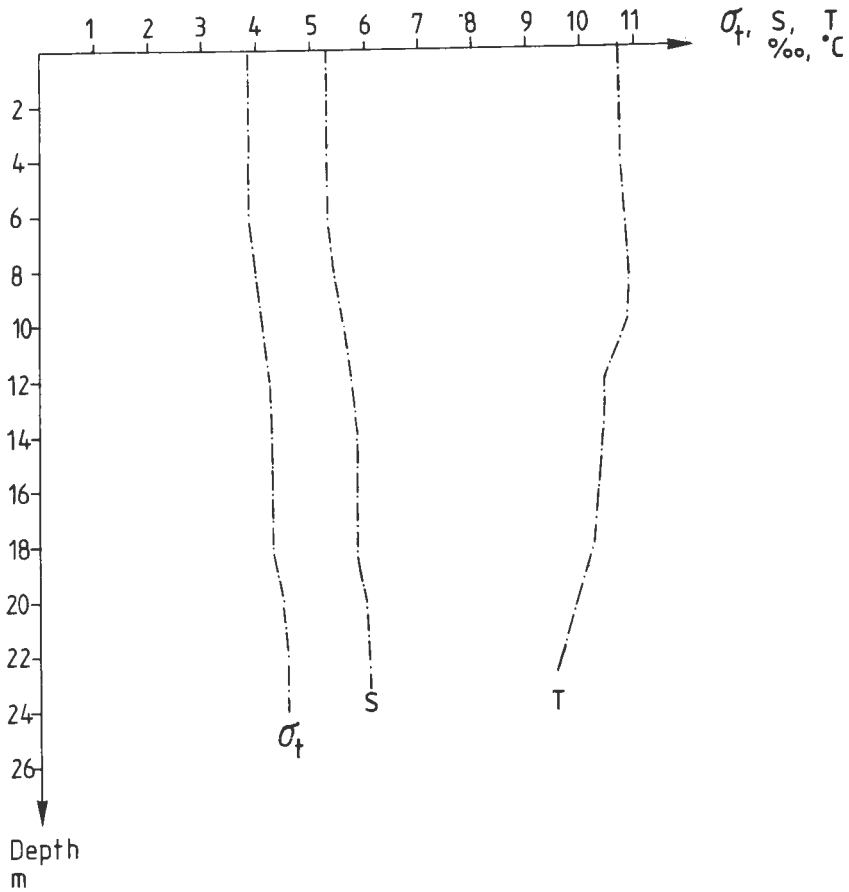


Figure 2. Vertical distribution of salinity (S), temperature (T) and density ( $\sigma_t$ ) 781002 time 17.45

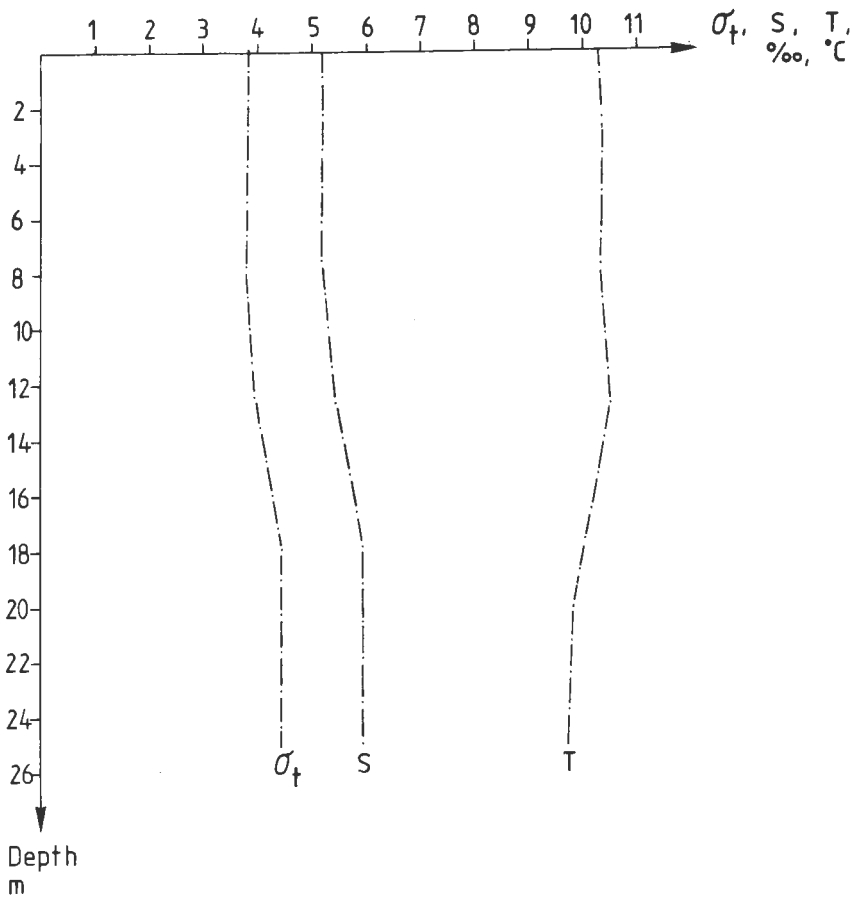


Figure 3. Vertical distribution of salinity (S), temperature (T) and density ( $\sigma_t$ ) 781003 time 15.00

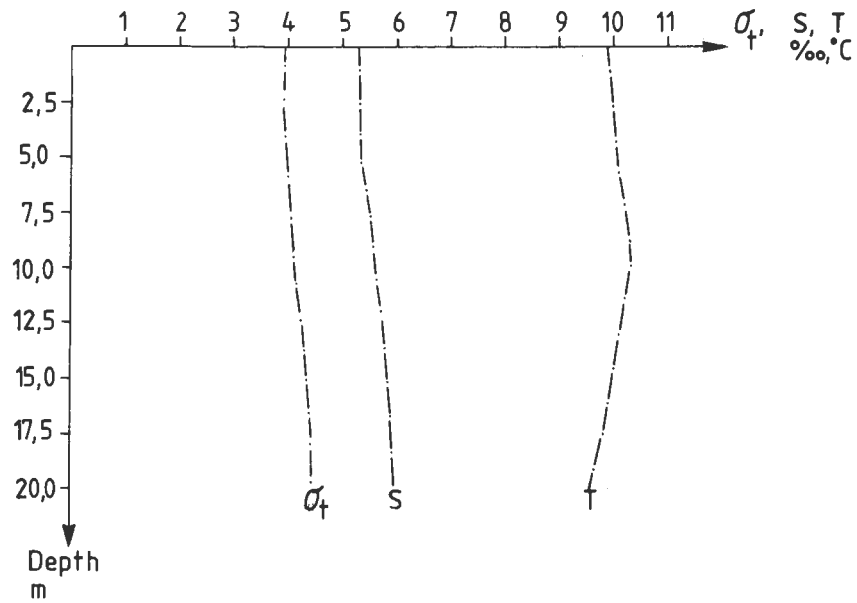


Figure 4. Vertical distribution of salinity (S), temperature (T) and density ( $\sigma_t$ ) 781005 time 6.00

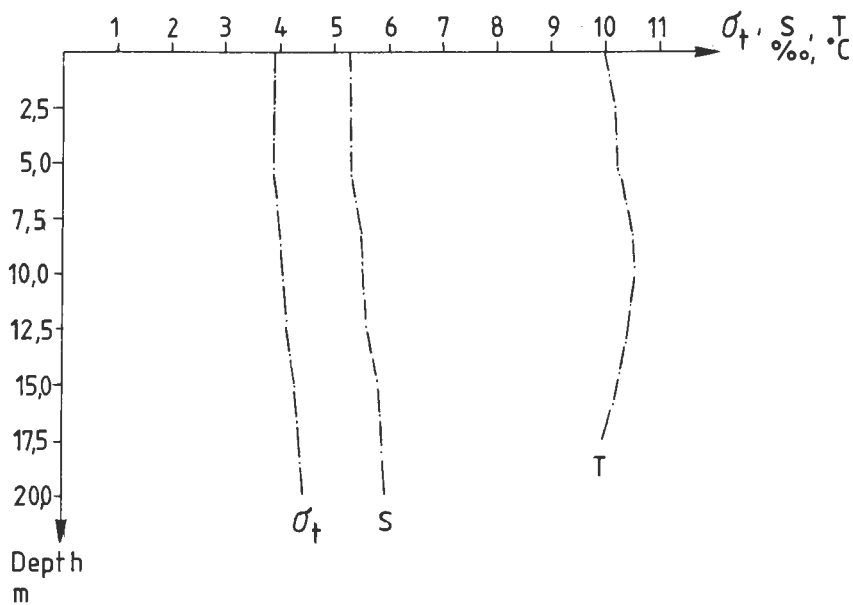


Figure 5. Vertical distribution of salinity (S), temperature (T) and density ( $\sigma_t$ ) 781005 time 8.15

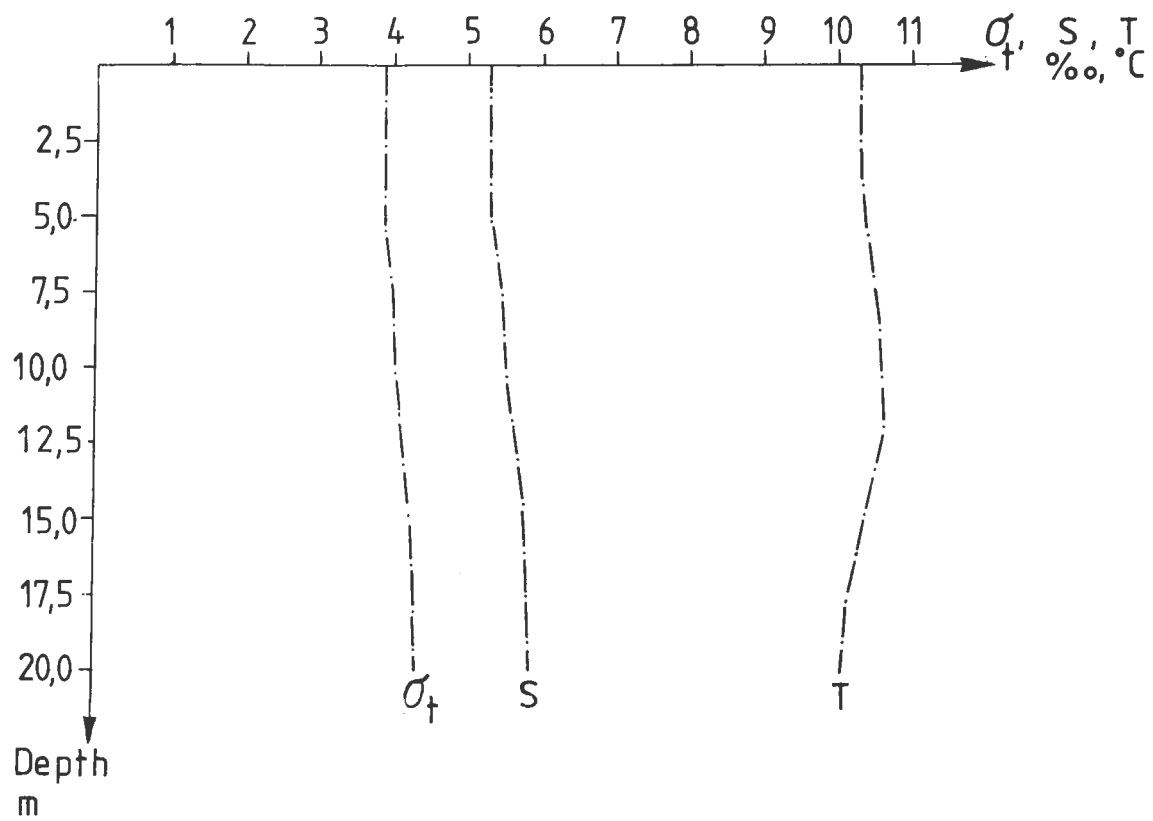


Figure 6. Vertical distribution of salinity (S), temperature (T) and density ( $\sigma_t$ ) 781005 time 8.30



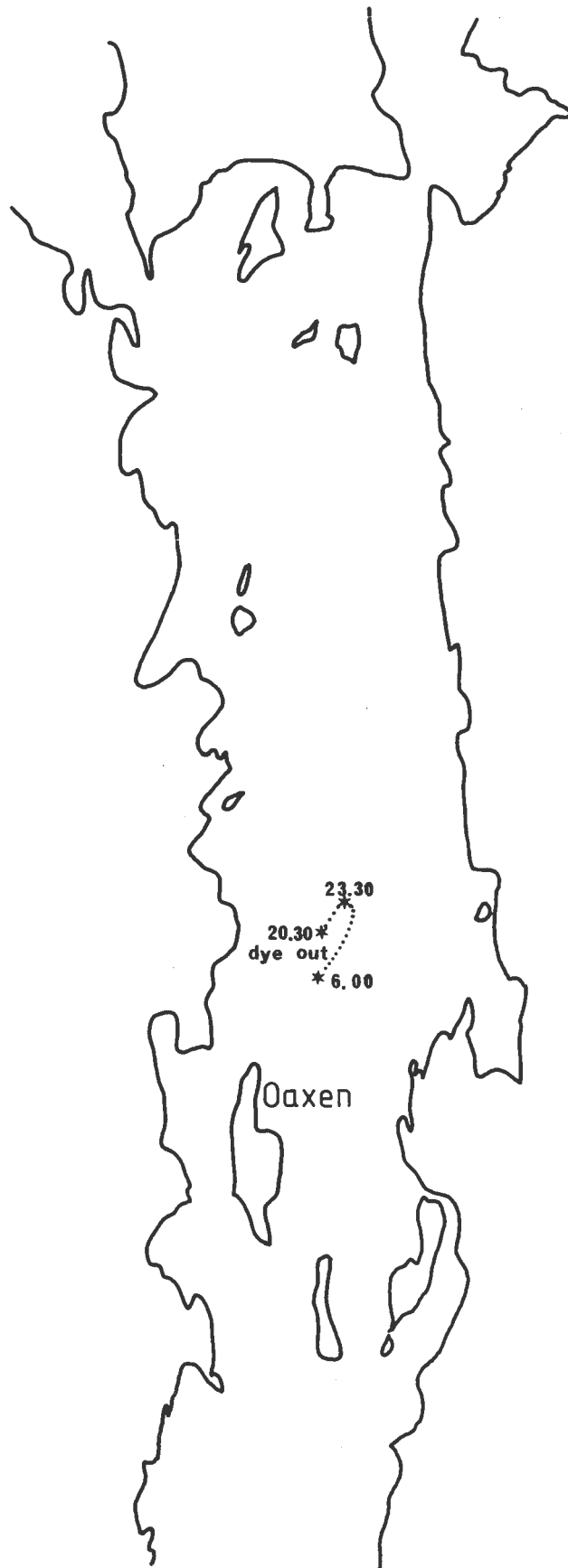


Figure 7. Track of the dye patch during experiment 1

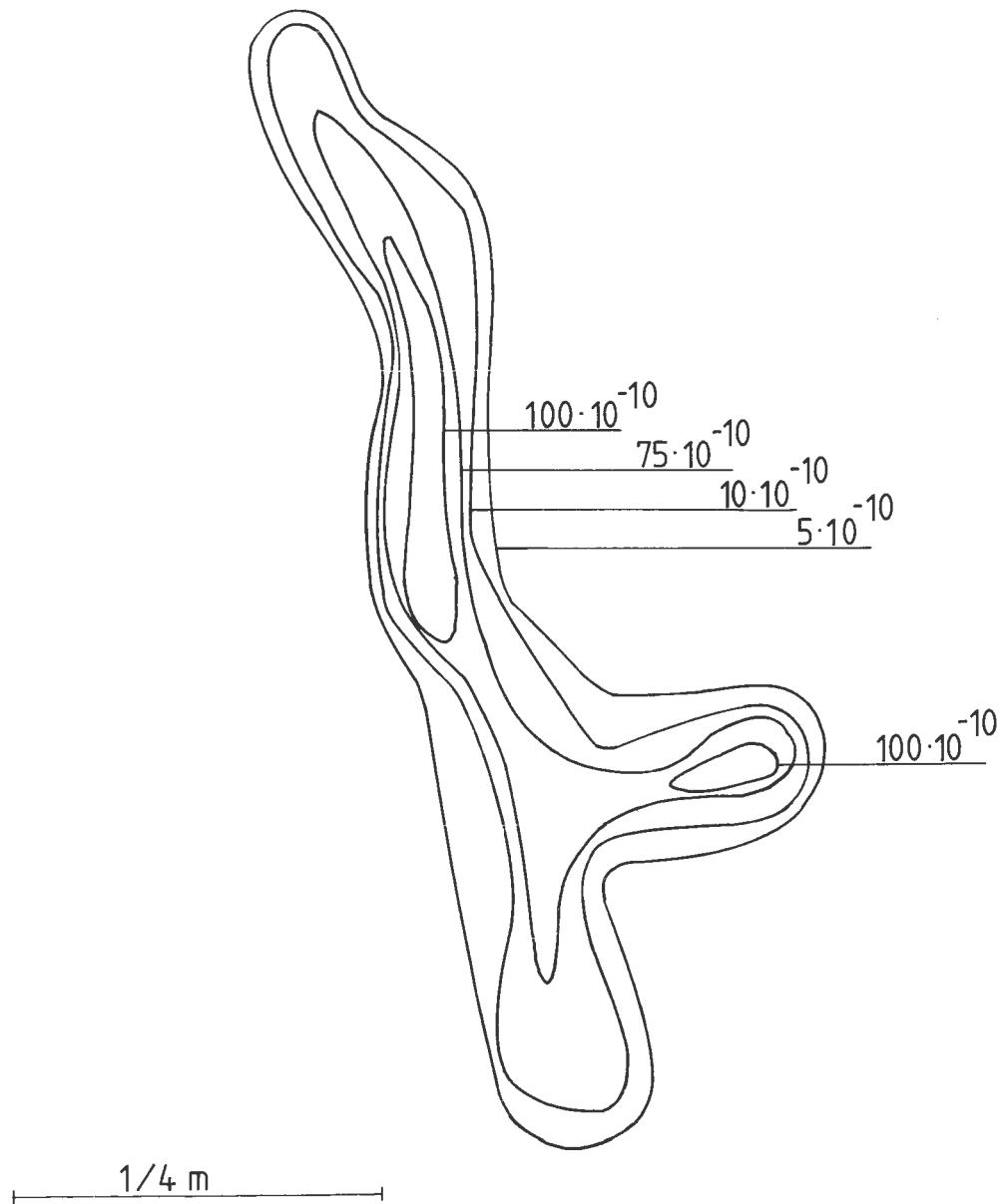


Figure 8. Dye distribution ( $t/m^3$ ). 781003 time 5.06-7.23.  
Experiment 1

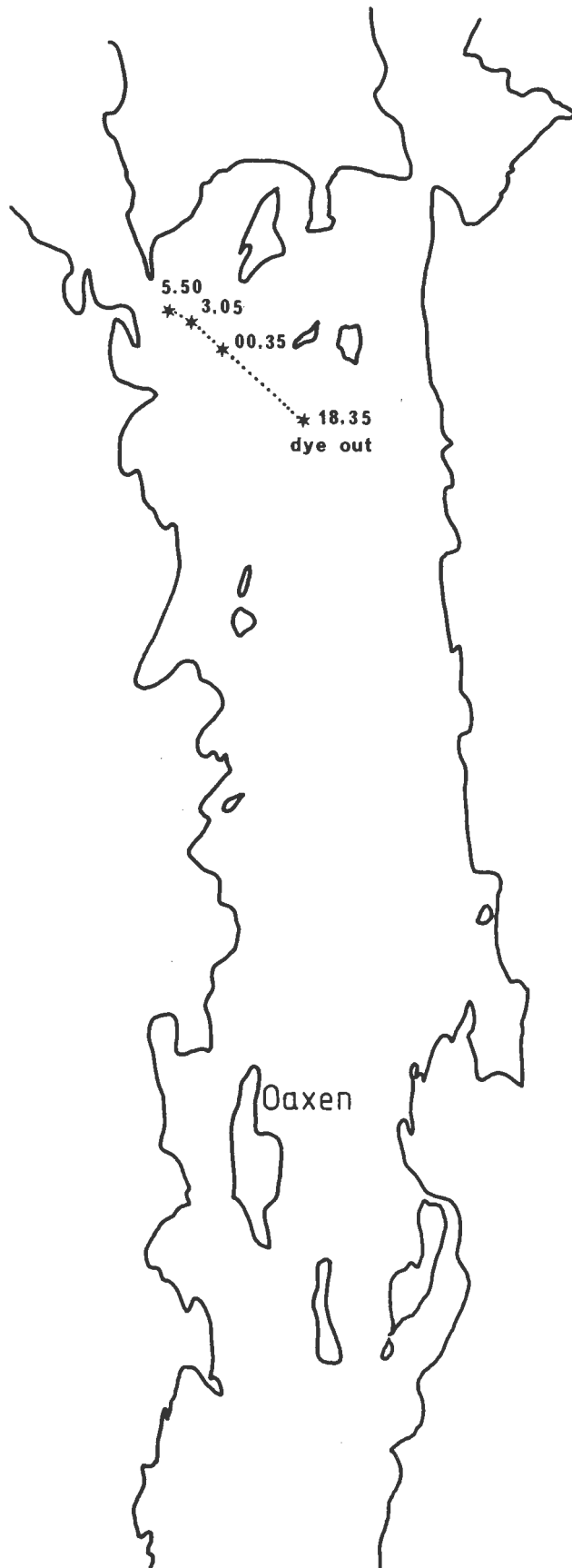


Figure 9. Track of the dye patch during experiment 2

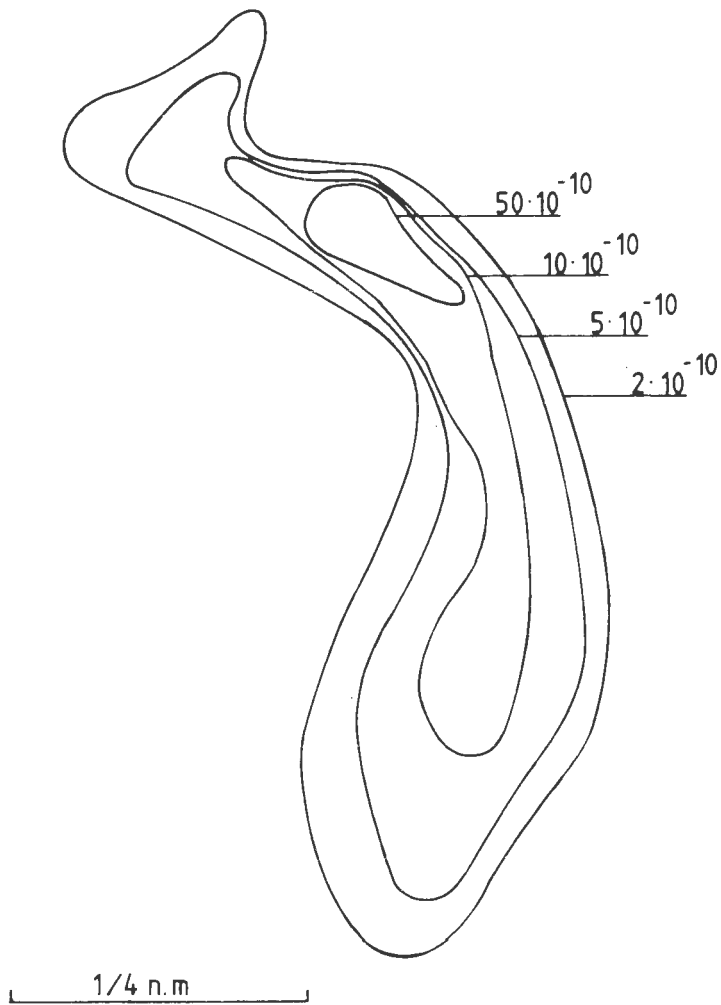


Figure 10. Dye distribution ( $\text{t/m}^3$ ). 781003 time 23.22-01.53.  
Experiment 2

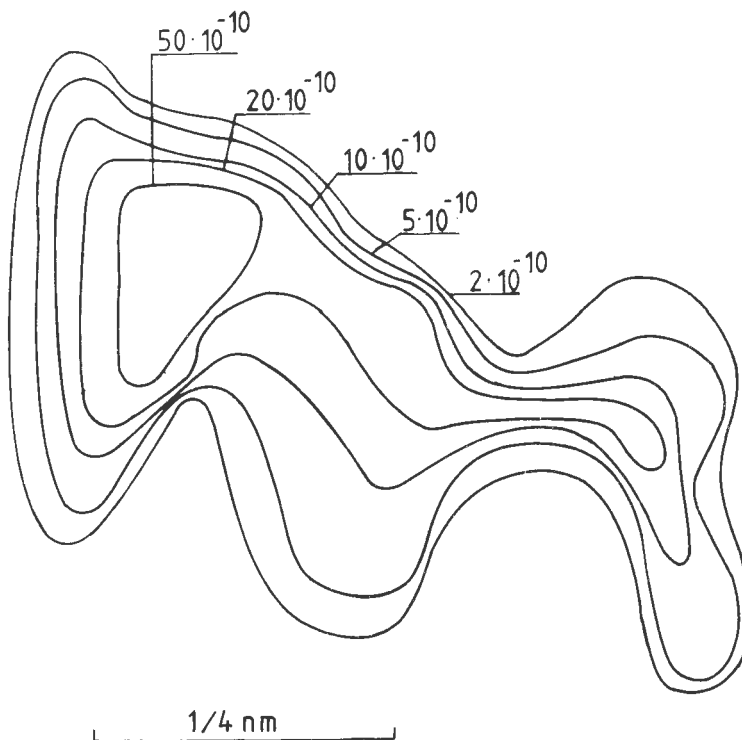


Figure 11. Dye distribution ( $\text{t/m}^3$ ). 781004 time 02.22-03.56  
Experiment 2

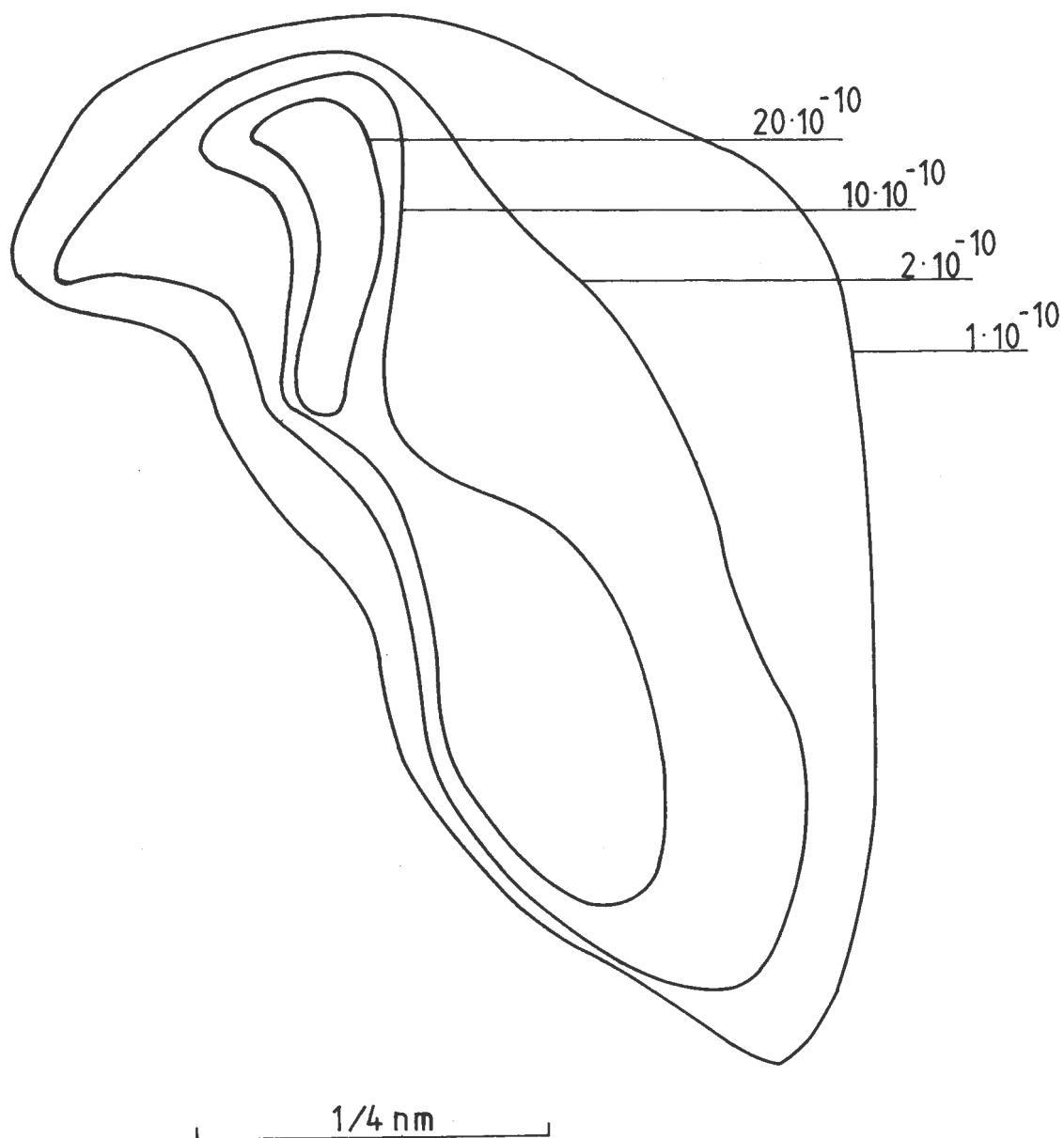


Figure 12. Dye distribution ( $t/m^3$ ). 781004 time 04.57-06.59.  
Experiment 2

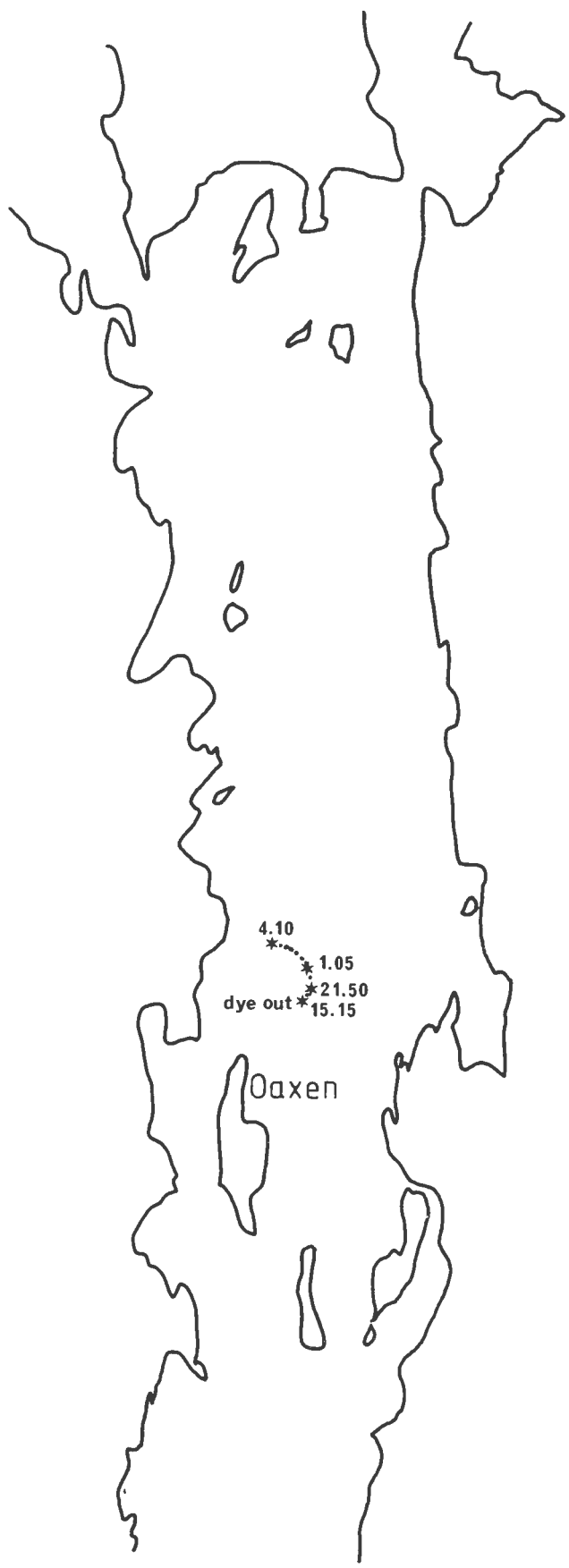


Figure 13. Track of the dye patch during experiment 3

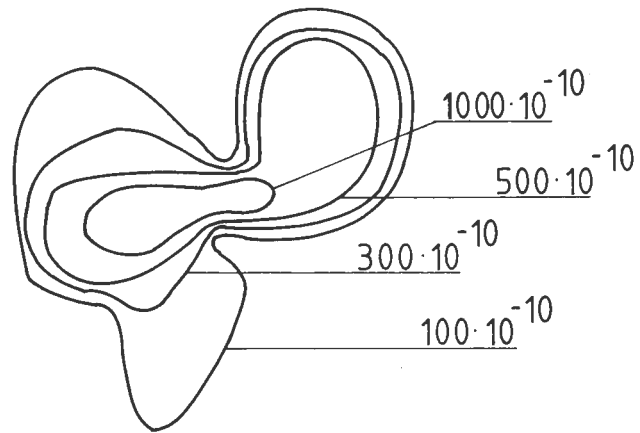


Figure 14. Dye distribution ( $\text{t/m}^3$ ). 781005 time 21.13-22.36.  
Experiment 3

1/4 nm

A horizontal scale bar with a bracket underneath, labeled "1/4 nm".

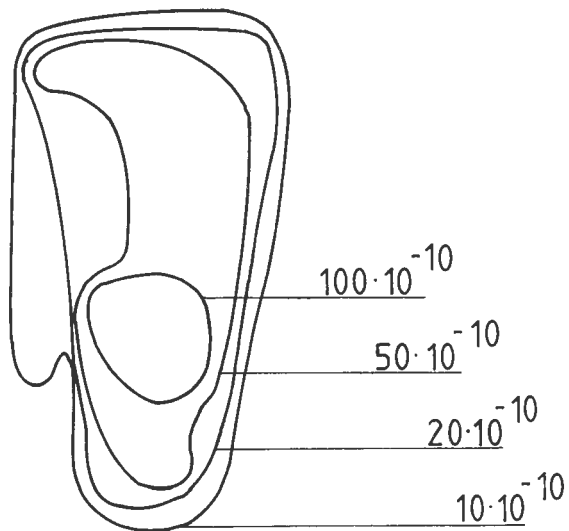


Figure 15. Dye distribution ( $\text{t/m}^3$ ). 781006 time 00.33-01.41.  
Experiment 3

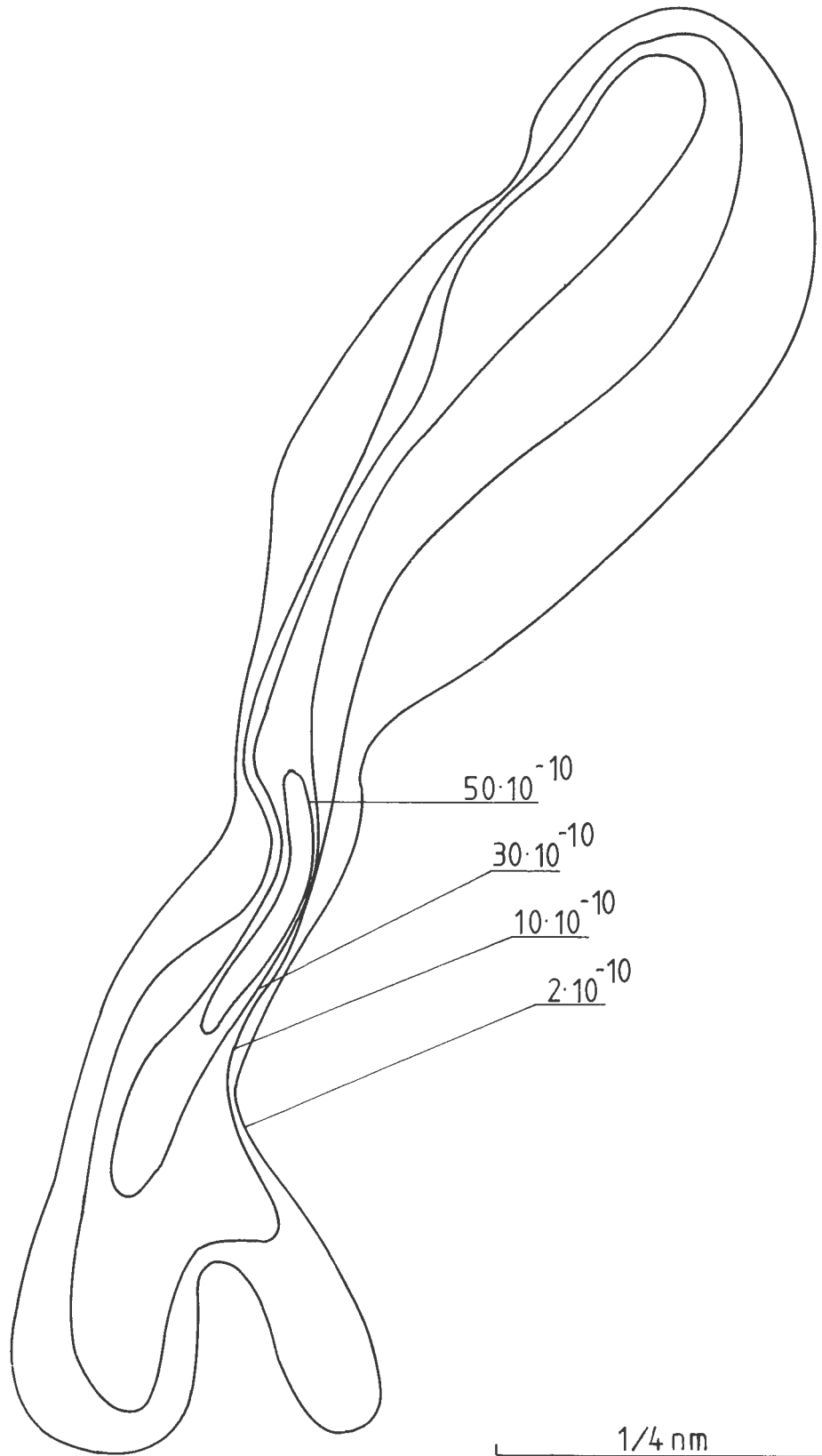


Figure 16. Dye distribution ( $\text{t/m}^3$ ). 781006 time 03.03-05.16.  
Experiment 3



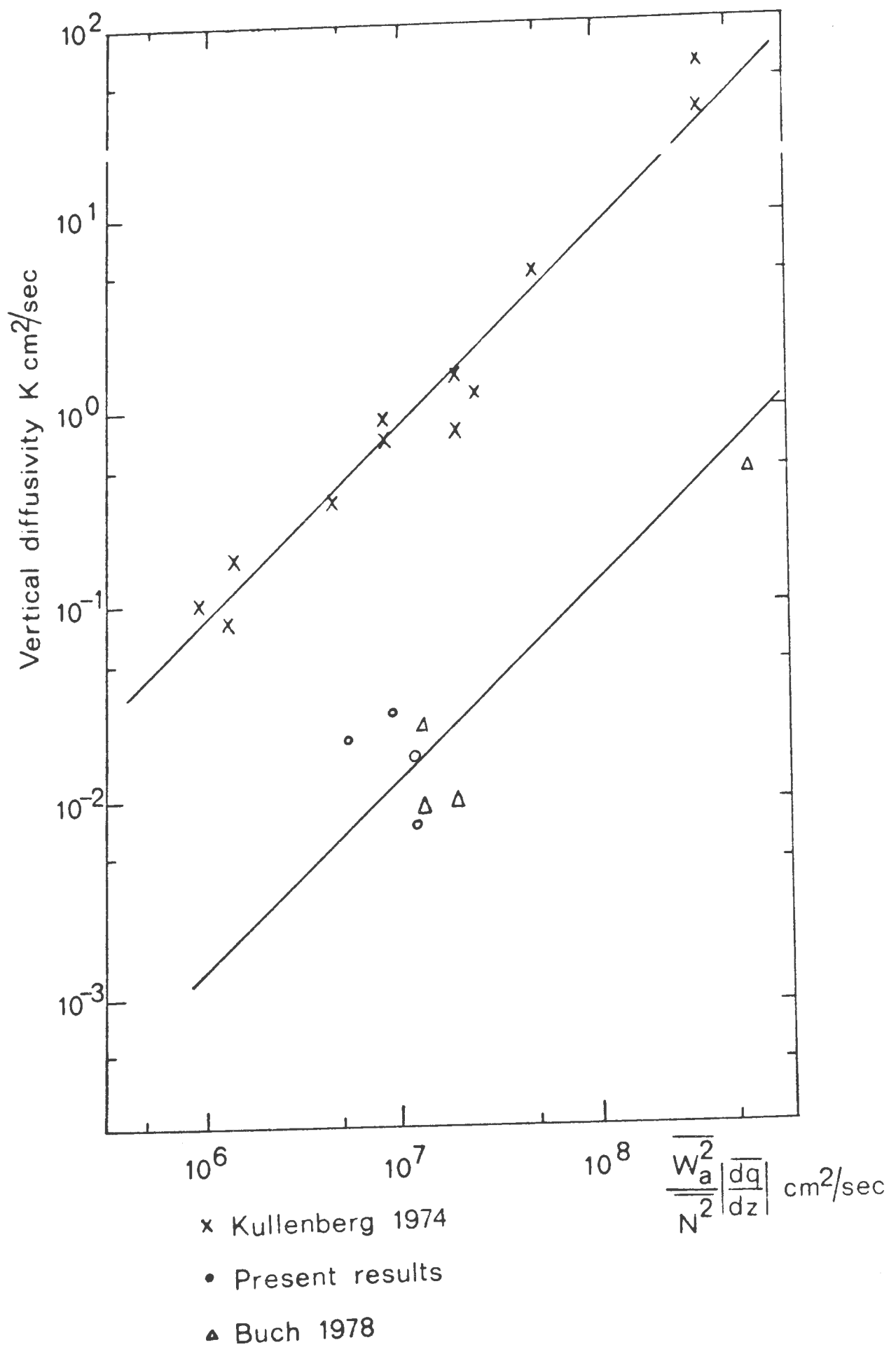
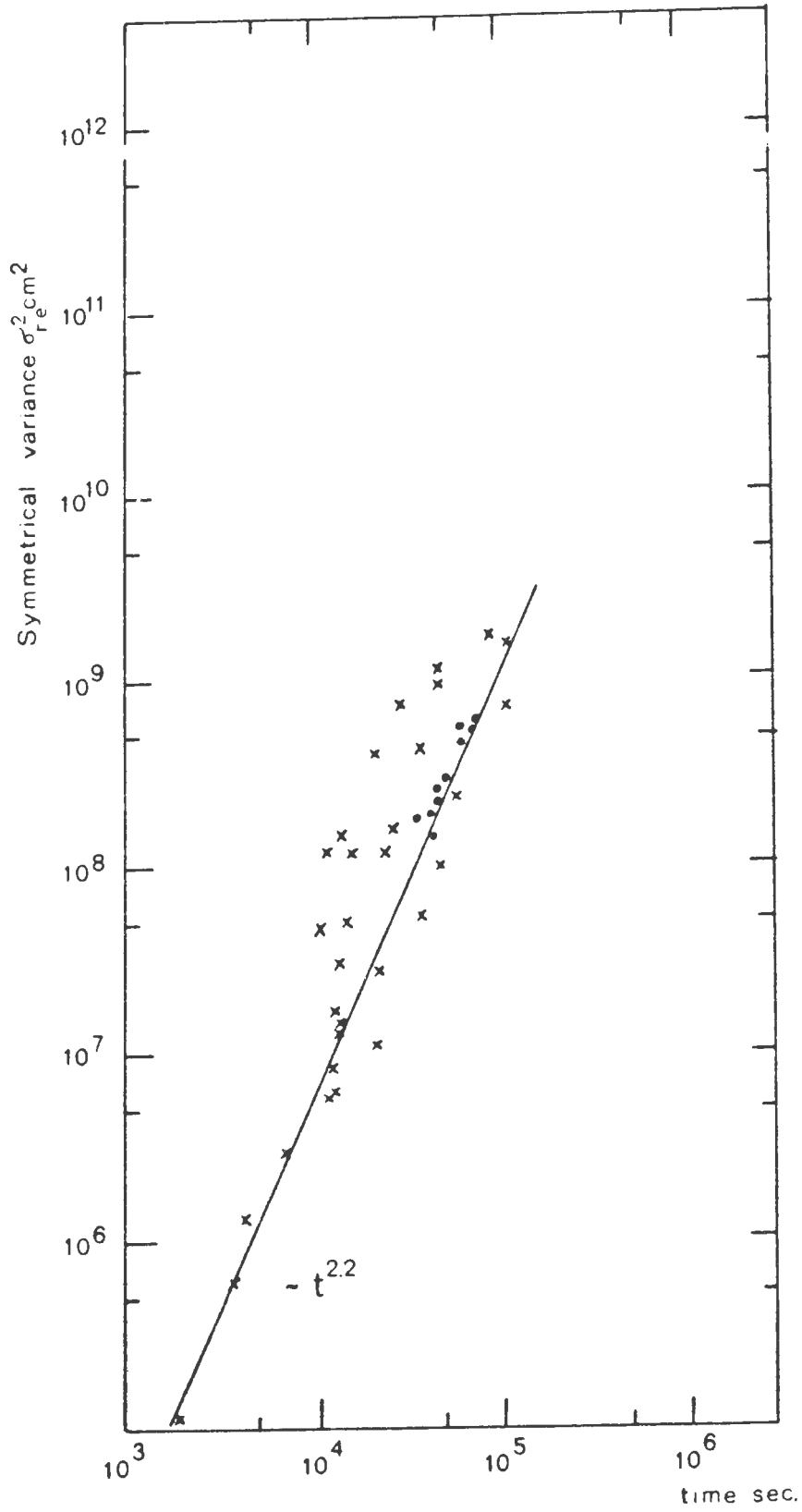


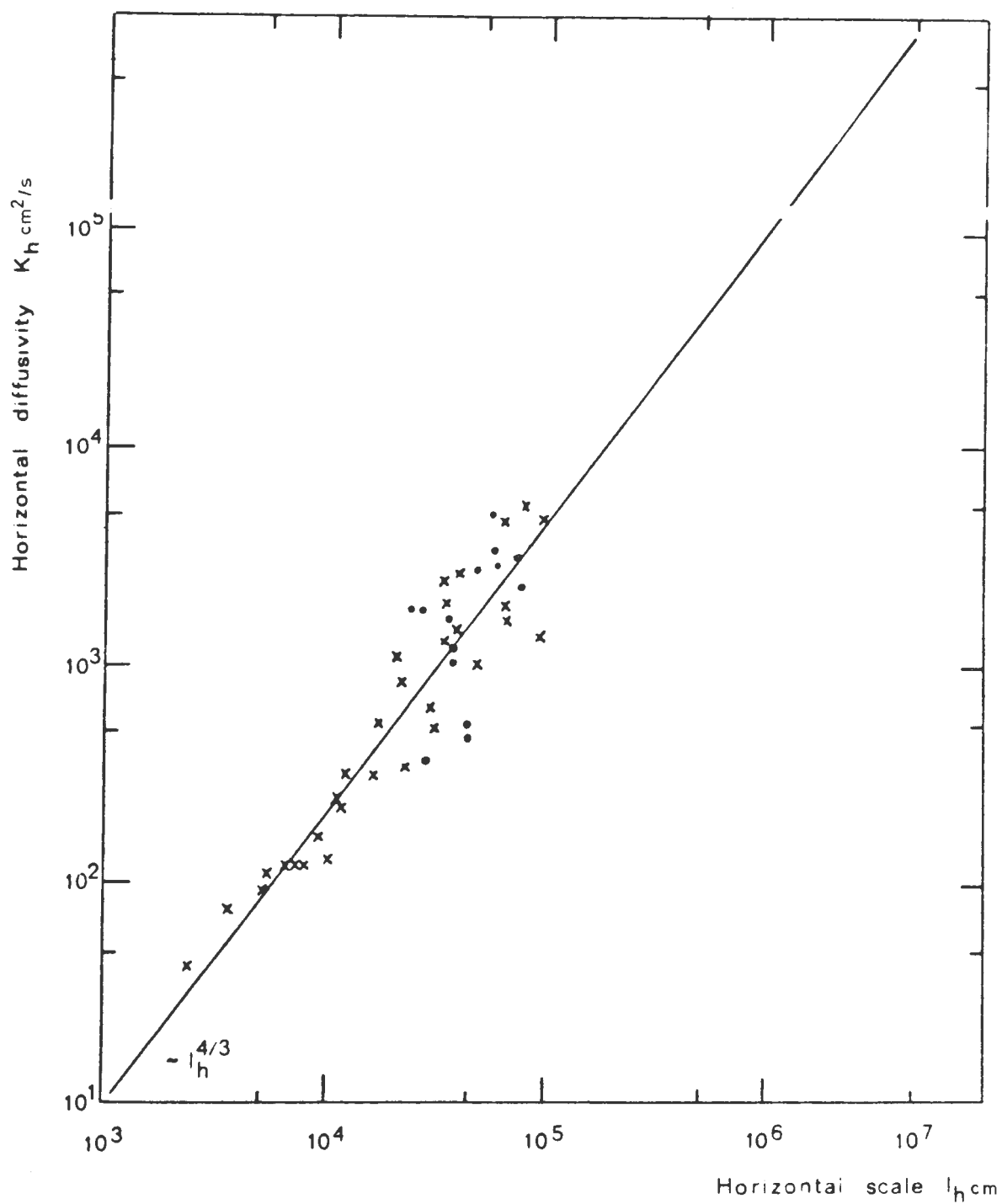
Figure 17. Vertical diffusion coefficient  $K$  versus the factor  $\frac{\overline{W_a^2}}{N^2} \left| \frac{dq}{dz} \right|$



× Kullenberg

• Present results

Figure 18. Symmetrical variance  $\sigma_{rc}^2$  versus  $t$



x Kullenberg 1974

• Present results

Figure 19. Horizontal diffusion coefficient  $K_h$  versus horizontal length scale  $l_h$

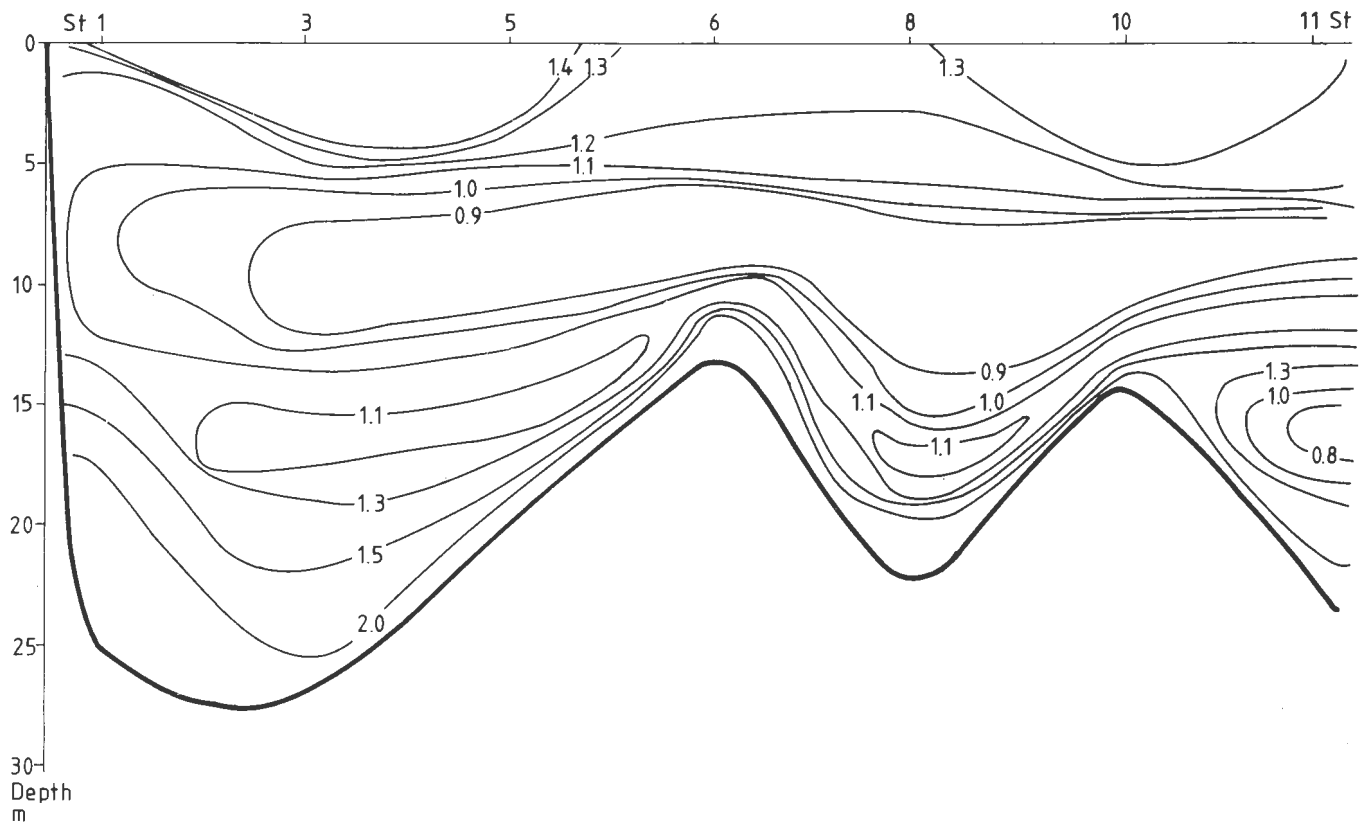


Figure 20. Distribution of  $c_p$ , attenuation coefficient due to particles. Longitudinal section. 781002

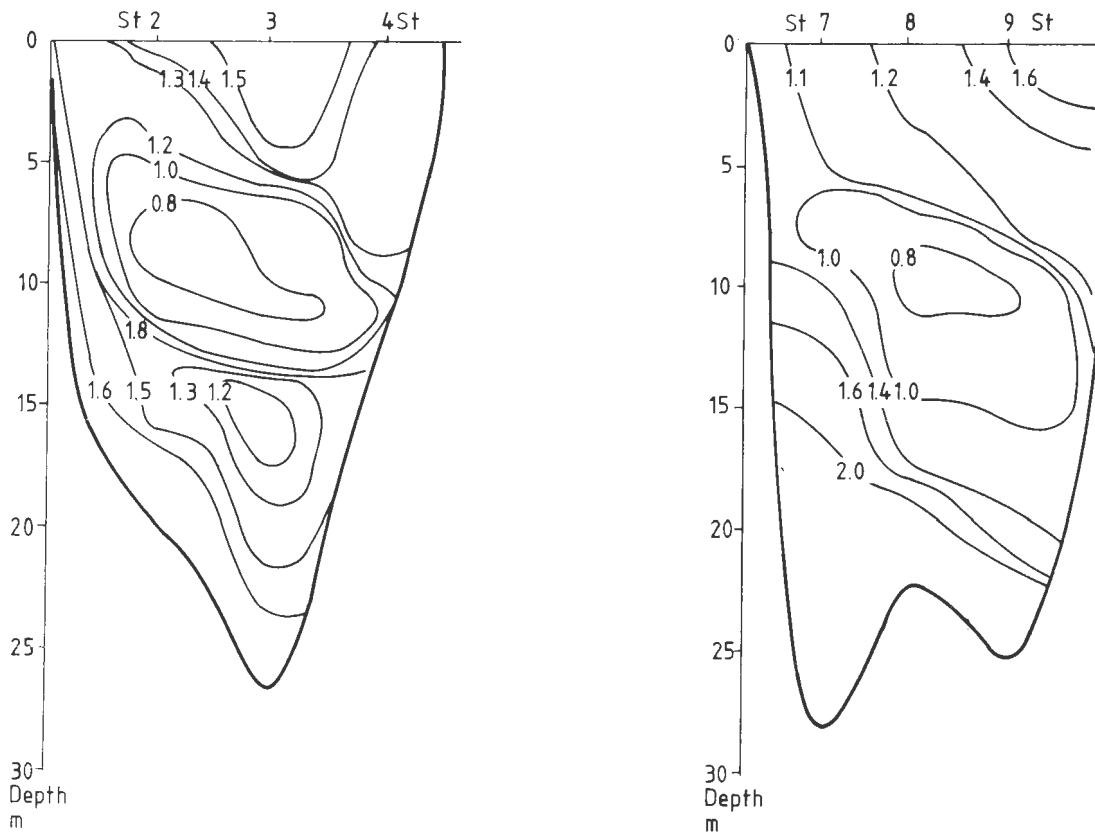


Figure 21. Distribution of  $c_p$ , attenuation coefficient due to particles. Transversal section. 781002

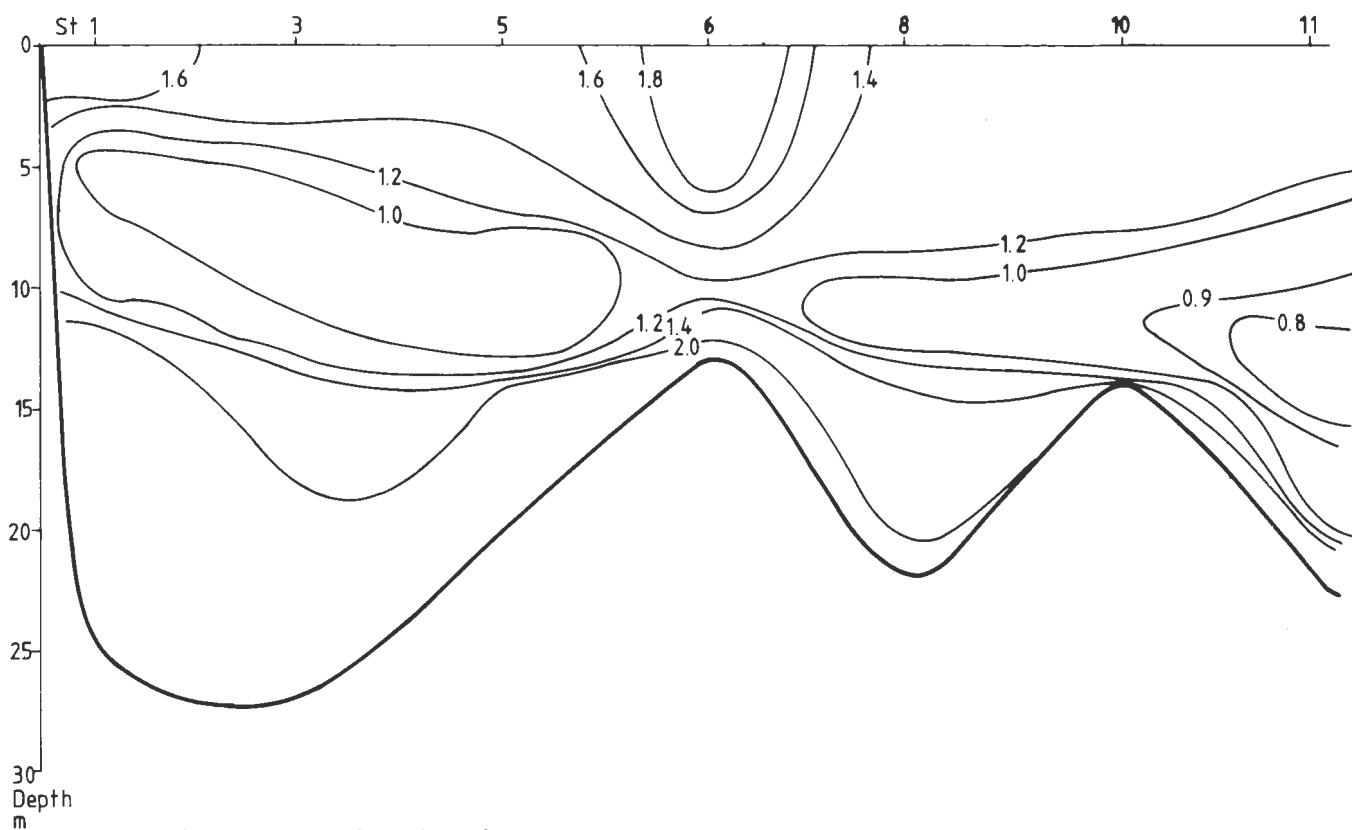


Figure 22. Distribution of  $c_p$ , attenuation coefficient due to particles. Longitudinal section. 781005

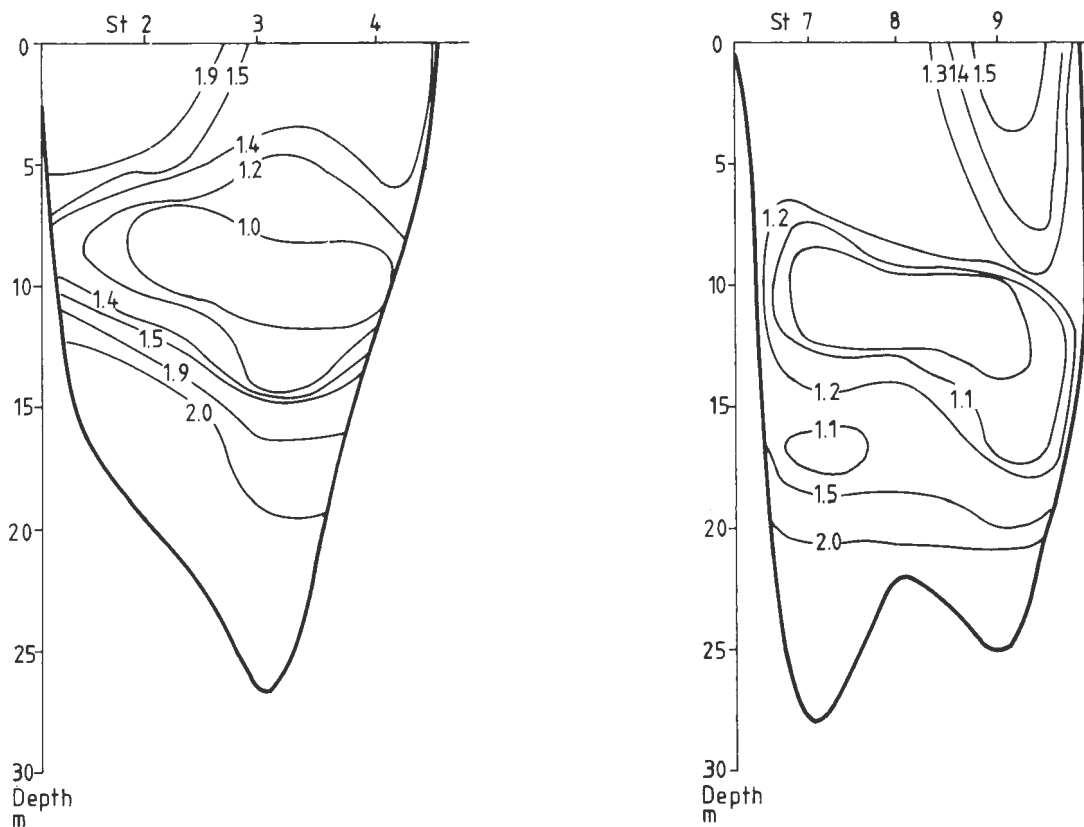


Figure 23. Distribution of  $c_p$ , attenuation coefficient due to particles. Transversal section. 781005

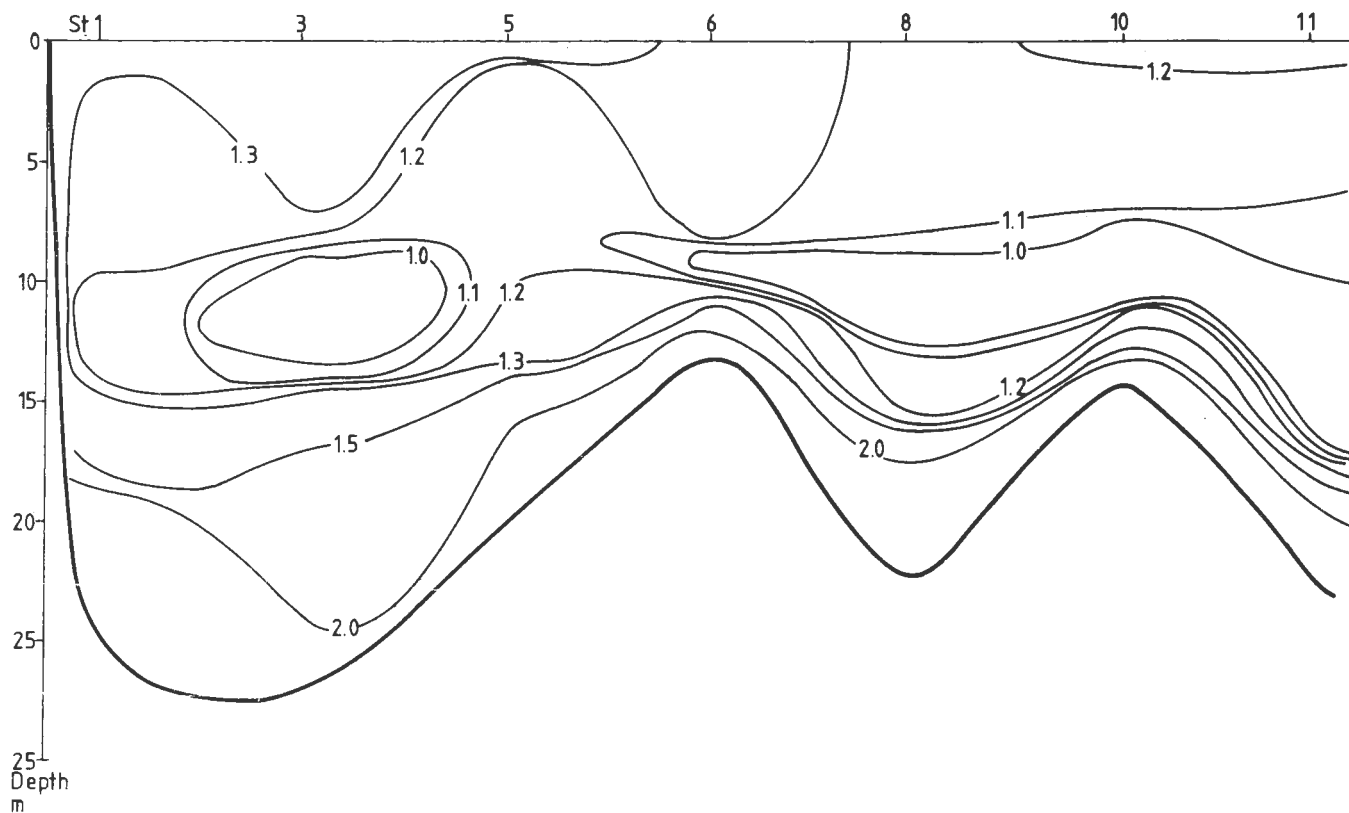


Figure 24. Distribution of  $c_p$ , attenuation coefficient due to particles. Longitudinal section. 781006

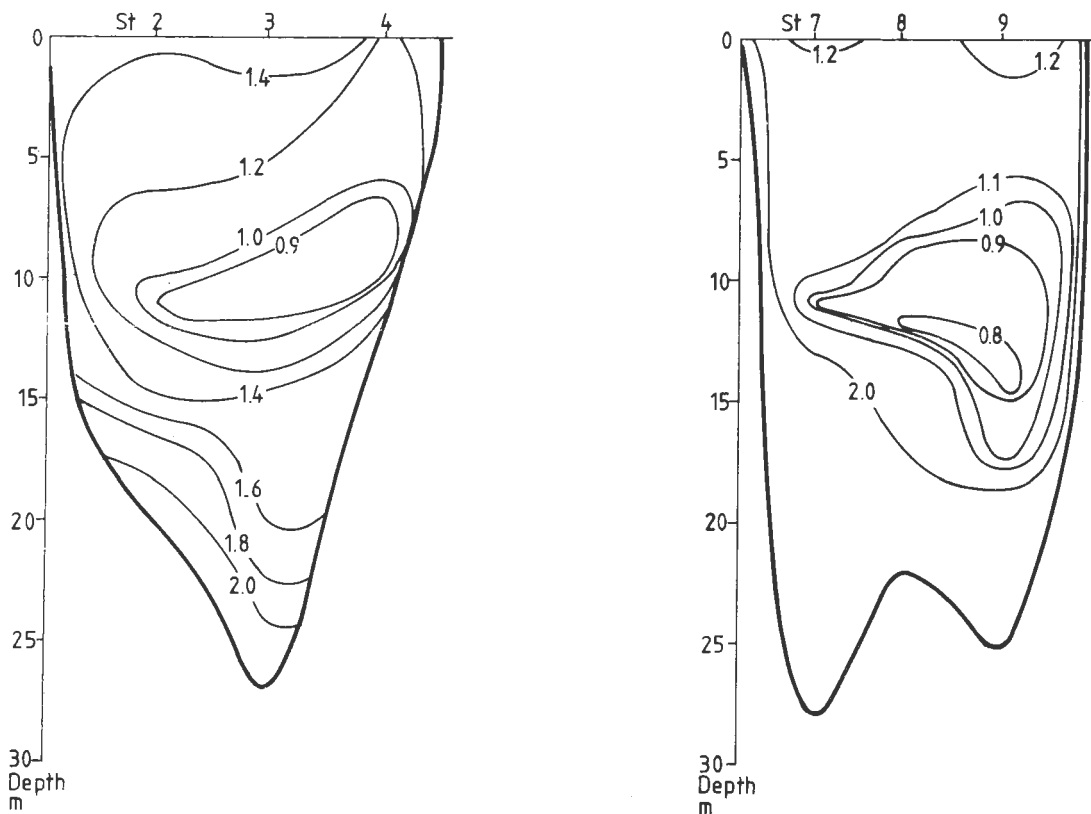


Figure 25. Distribution of  $c_p$ , attenuation coefficient due to particles. Transversal section. 781006

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SVERIGES METEOROLOGISKA OCH HYDROLOGISKA INSTITUT  
Folkborgsvägen 1, Box 923, 601 19 Norrköping. Telefon 011-108000.